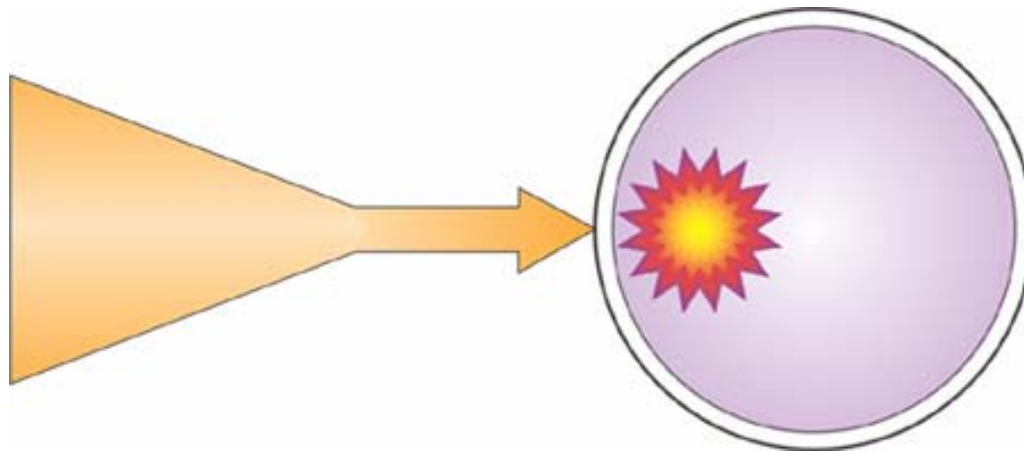
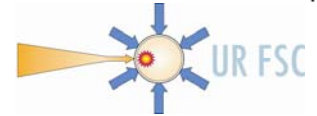


# FUEL ASSEMBLY: THEORY AND EXPERIMENTS

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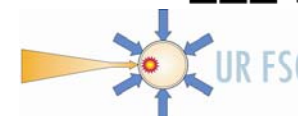
R. Betti, C. Zhou, A. Solodov, W. Theobald,  
K. Anderson, C. Stoeckl, D. Meyerhofer  
*Fusion Science Center, LLE*  
C. Li and R. Petrasso  
*Fusion Science Center, MIT*

4<sup>th</sup> Fusion Science Center Meeting  
Livermore, CA  
28-29 August 2006

# Outline

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- **Review of the hydrodynamic theory for fast ignition**
- **Two dimensional simulations of FI fuel assembly**
- **Gain curves for fast ignition**
- **OMEGA experiments for FI fuel assembly on CH targets**
- **Cryogenic target design for fuel assembly on OMEGA**

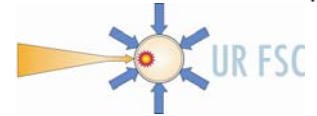
**High areal densities ( $\rho R$ ) and low-implosion velocities ( $V_i$ ) lead to high energy gains (assuming that ignition occurs)**

$$\mathbf{G} = \frac{\mathbf{E}_{Fusion}}{\mathbf{E}_{Laser}} \sim \frac{\theta}{V_i^{1.25}}$$

$$\theta = \frac{1}{1 + 7/\rho R} = \text{burned fraction}$$

- Higher  $\rho R \rightarrow$  faster burn and longer burn time
- Lower  $V_i \rightarrow$  more fuel mass for the same kinetic/laser energy

# The hydrodynamics of fast ignition depends on three parameters: gain, density, and areal density



$$Gain^a \sim V_i^{-1.25} (1 + 7/\rho R)^{-1} \Rightarrow \frac{743}{1 + 30 / E_L^{1/3} (kJ)}$$

$$\rho R^a \sim E_L^{0.33} / \alpha^{0.55}$$

$$\rho^a \sim V_i / \alpha$$

**Fast ignition implosion:**

- Low velocity  $V_i$
- Low adiabat  $\alpha$
- Large mass

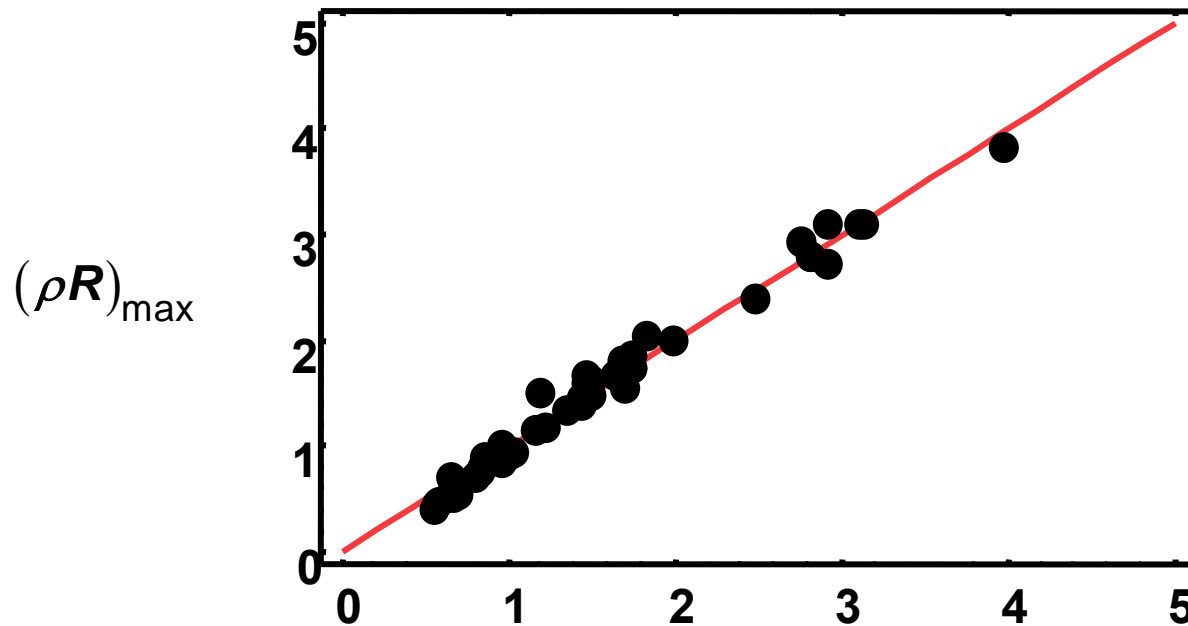
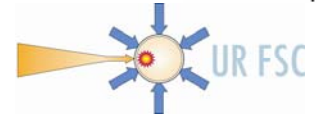
$$E_{ig}^b (kJ) \approx 11 \left( \frac{400}{\rho_{g/cc}} \right)^{1.85}$$

High  $\rho$  is required for fast ignition

<sup>a</sup> R. Betti and C. Zhou, Phys. Plasmas 12, 110702 (2005)

<sup>b</sup> S. Atzeni, Phys. Plasmas 6, 3316 (1999)

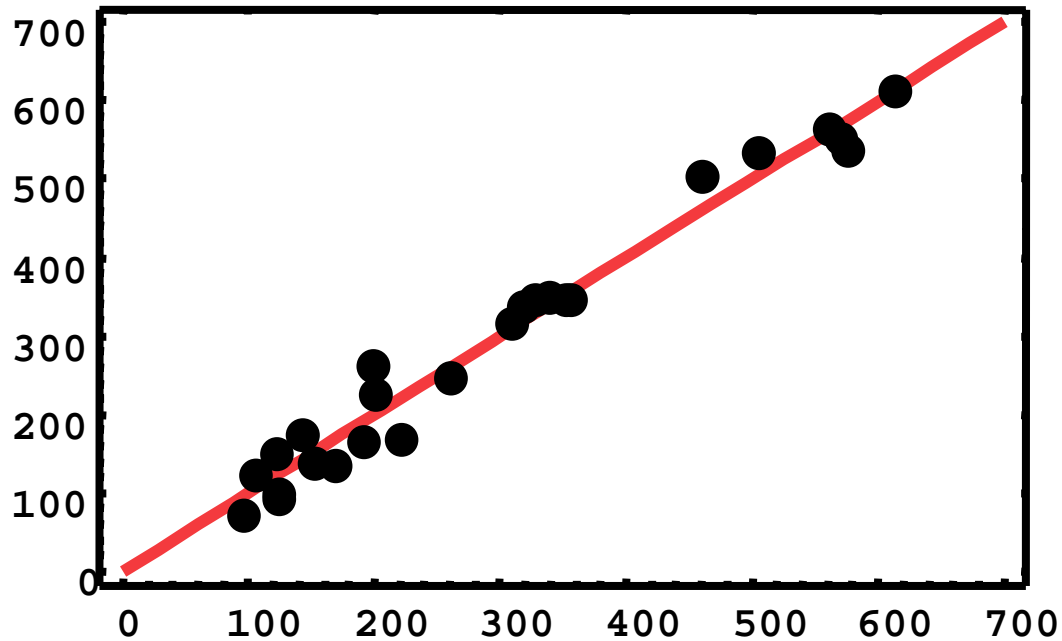
**We have increased the implosion simulation database.  
The areal density increases for lower adiabats  
and greater energies. It is independent of velocity.**



$$(\rho R)_{\max}^{fit} = \frac{1.2}{\alpha^{0.57}} \left( \frac{E_L (kJ)}{100} \right)^{0.35} \left( \frac{V_i (cm / s)}{3 \cdot 10^7} \right)^{0.1}$$

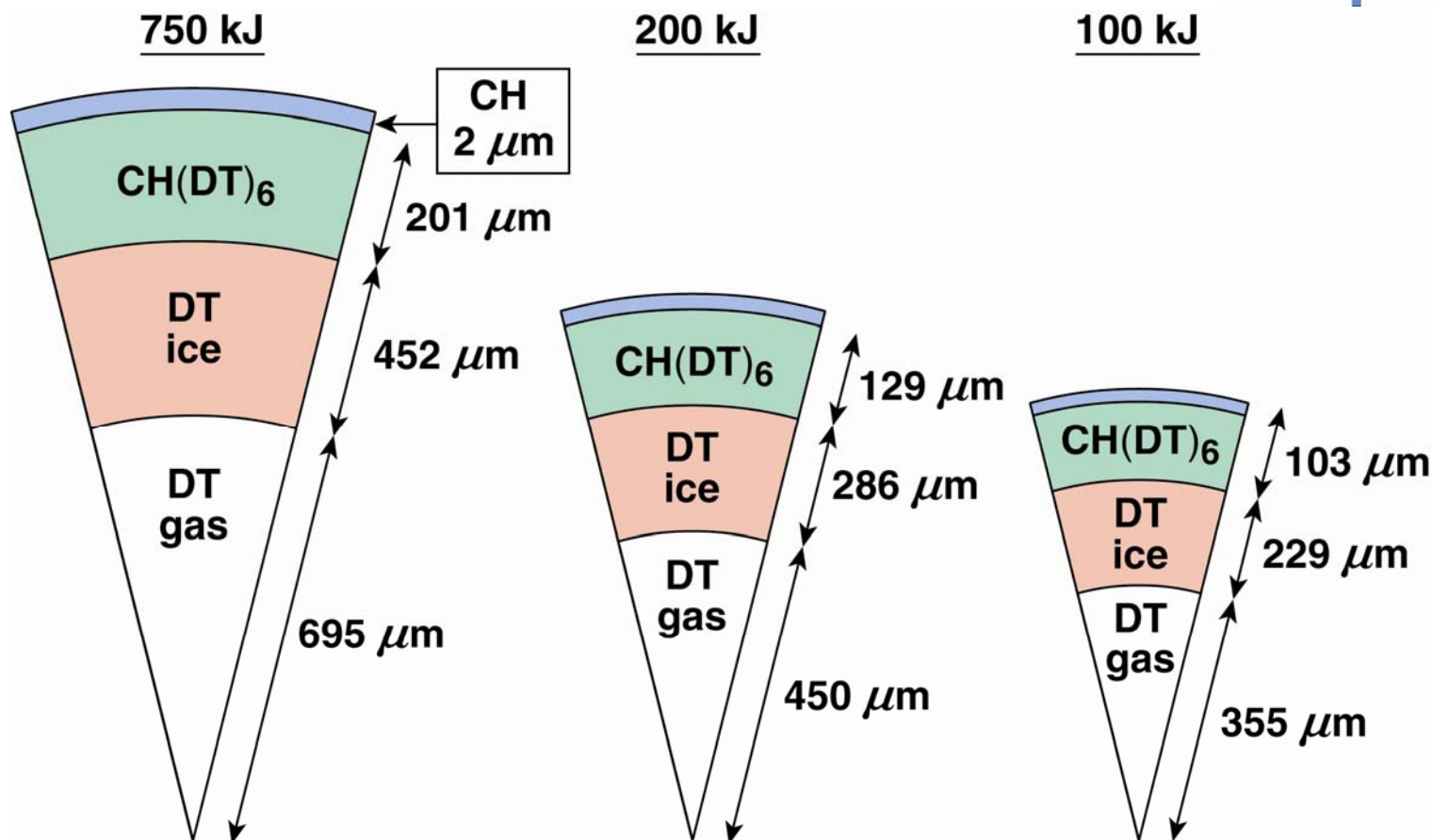
# The density increases with high velocities and lower adiabats

$$\langle \rho \rangle_{\rho R}$$

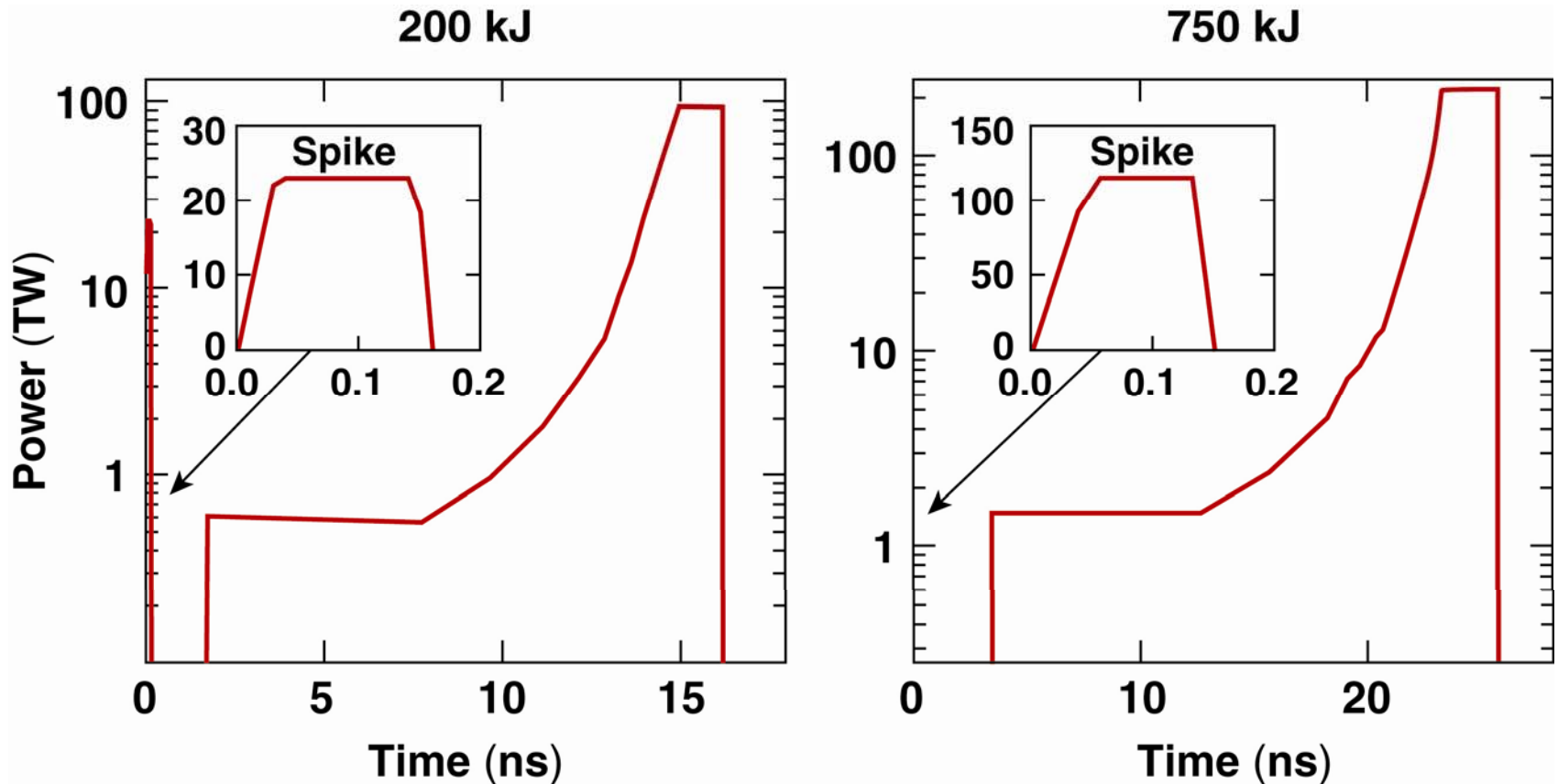


$$\langle \rho \rangle_{\rho R}^{fit} = \frac{440}{\alpha} \left( \frac{V_i (\text{cm} / \text{s})}{3 \cdot 10^7} \right)^{0.93}$$

# Optimized fast-ignition cryo-targets are thick shells of wetted foam with initial aspect ratio $\sim 2$

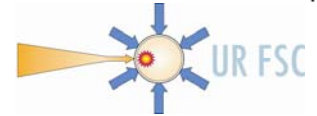


# Fast-ignition targets require long laser pulses and high contrast ratios ( $\sim 100$ to $150$ ) within the capabilities of the NIF





# Optimized fast-ignition implosions are NOT affected by the Rayleigh–Taylor instability



- Number of e-folds for the most dangerous RT modes

$$Ne_{k\Delta=1} \approx \frac{6V_i(\text{cm/s})}{3 \cdot 10^7} \left[ \frac{1}{\langle \alpha \rangle^{0.3}} - 0.08 \left( \frac{\alpha_{\text{outer}}}{\langle \alpha \rangle} \right)^{0.6} \right]$$

$$V_i \approx 1.7 \cdot 10^7$$

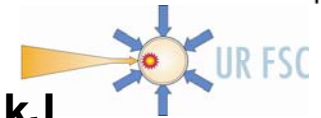
$$\alpha_{\text{inner}} \approx 0.7$$

$$\alpha_{\text{outer}} \approx 2$$

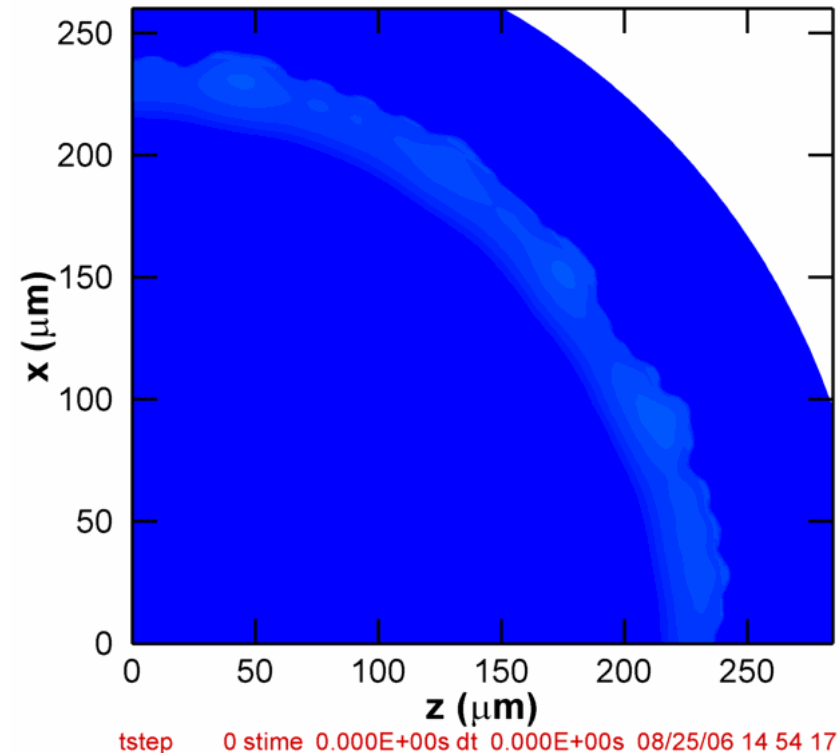
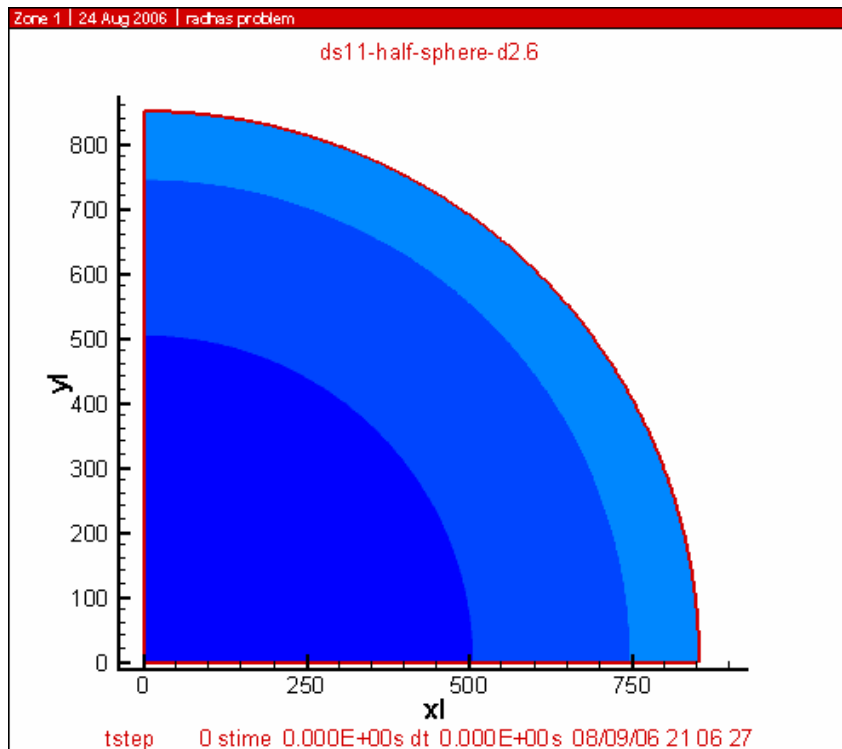


$$Ne_{k\Delta=1} \approx 3$$

# 2-D multimode simulations of a 250 kJ fast-ignition target implosion show little RT growth



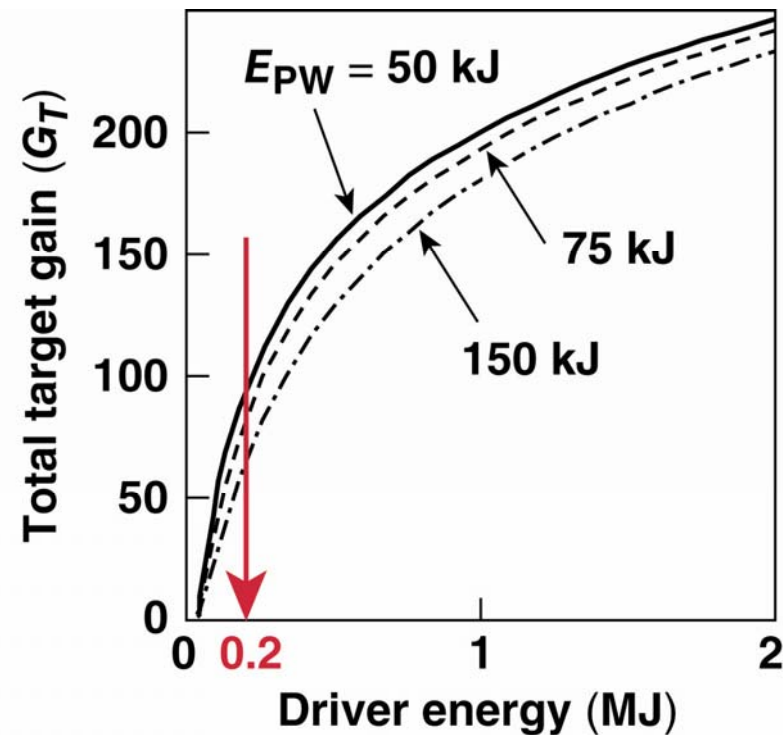
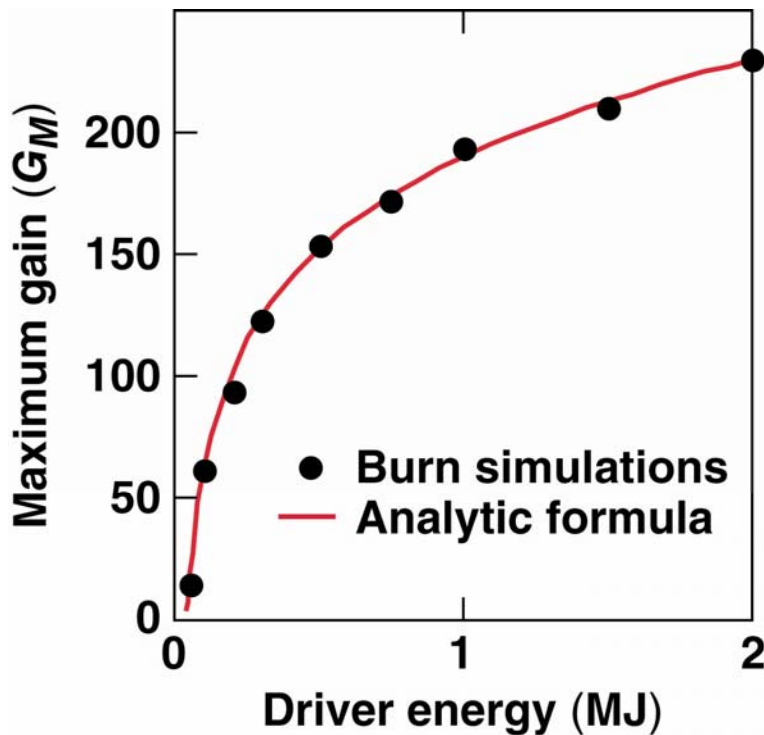
Laser imprinting, Modes  $\ell = 4$  to 100, 1-THZ SSD,  $E_L = 250$  kJ,  
 $V = 2 \cdot 10^7$  cm/s,  $\alpha = 0.7$ , IFAR = 21



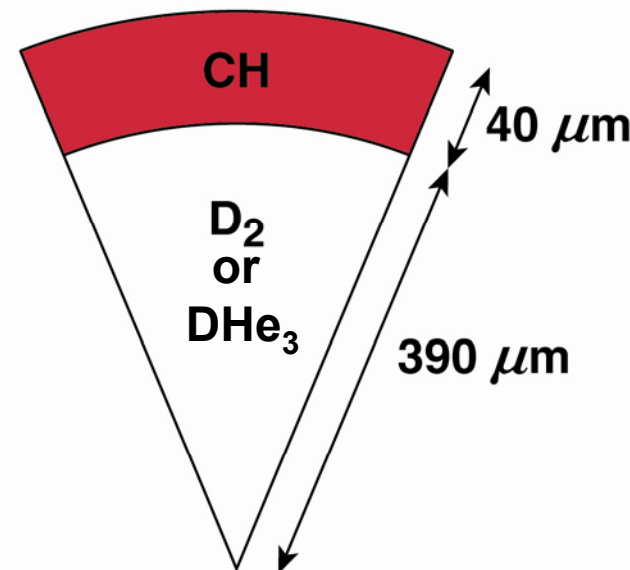
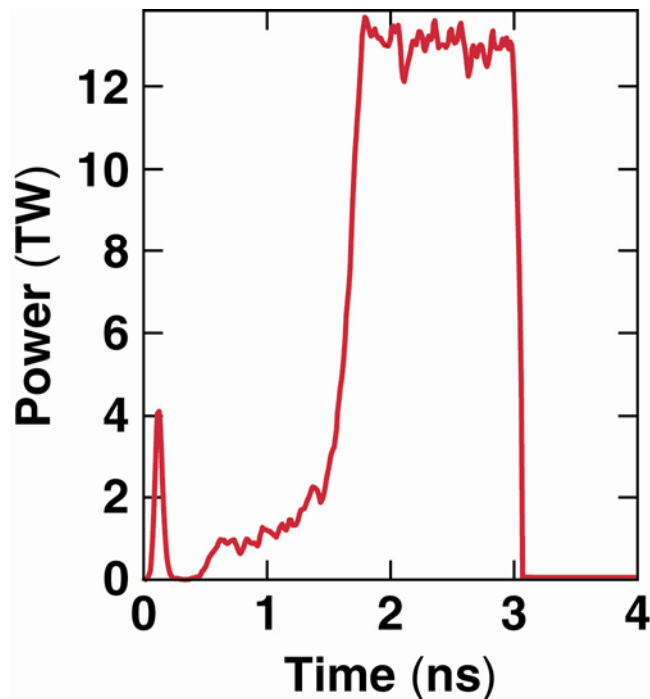
# High gains are possible with small drivers with energy as low as 200 kJ

$$G_M = \frac{E_{\text{Fusion}}}{E_{\text{Driver}}}$$

$$G_T = \frac{E_{\text{Fusion}}}{E_{\text{Driver}} + E_{\text{Petawatt}}}$$

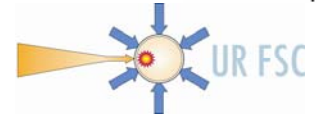


# Low-adiabat, low- $V_i$ implosions of CH targets are used to study fast-ignition fuel assemblies

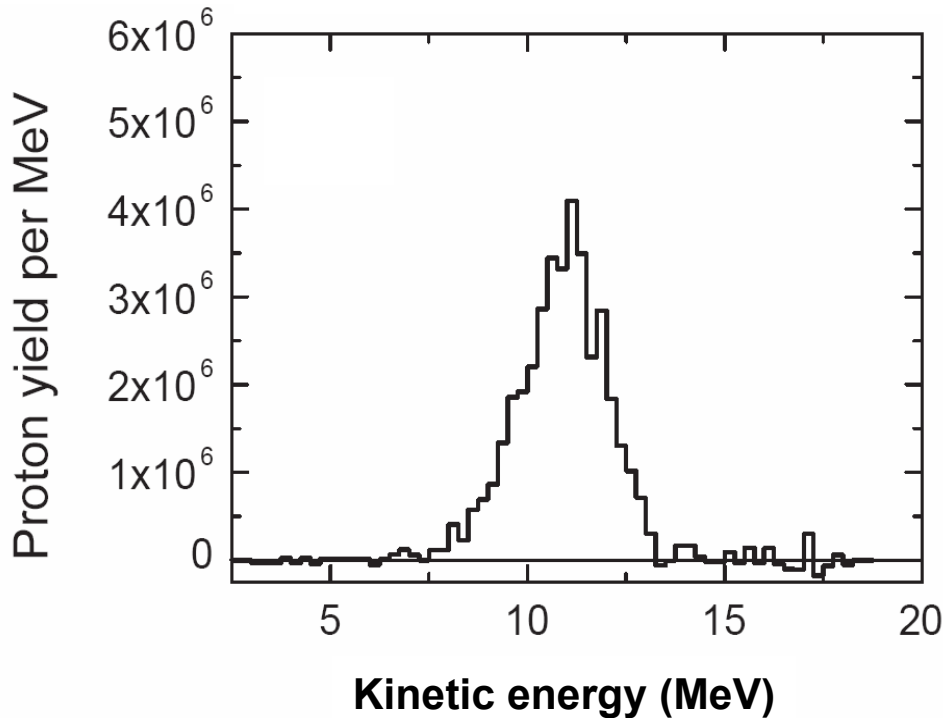


$$\alpha \approx 1.3, V_i \approx 2 \cdot 10^7 \text{ cm/s}$$

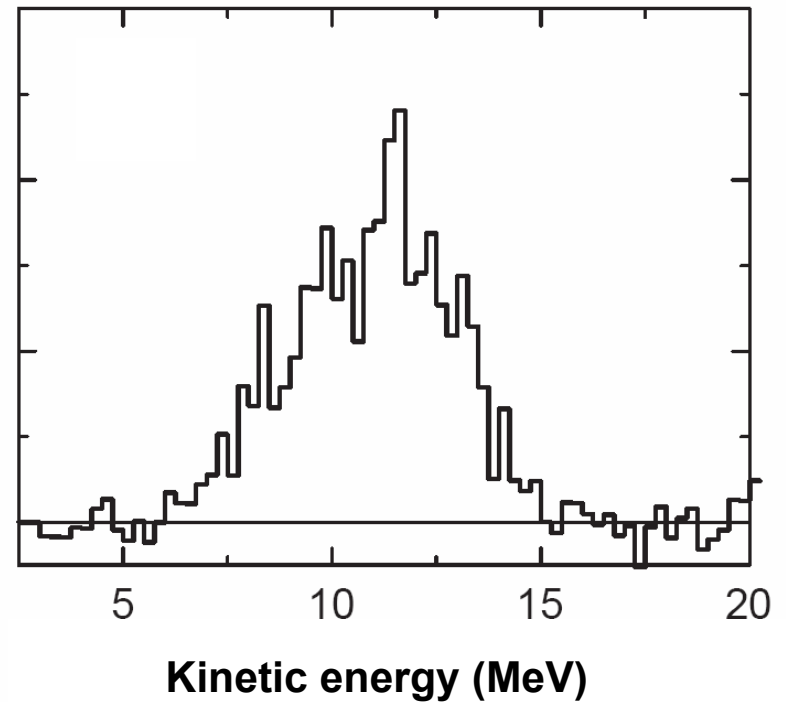
# Well defined downshifted proton spectra are measured for both D2 and DHe3 gas fills



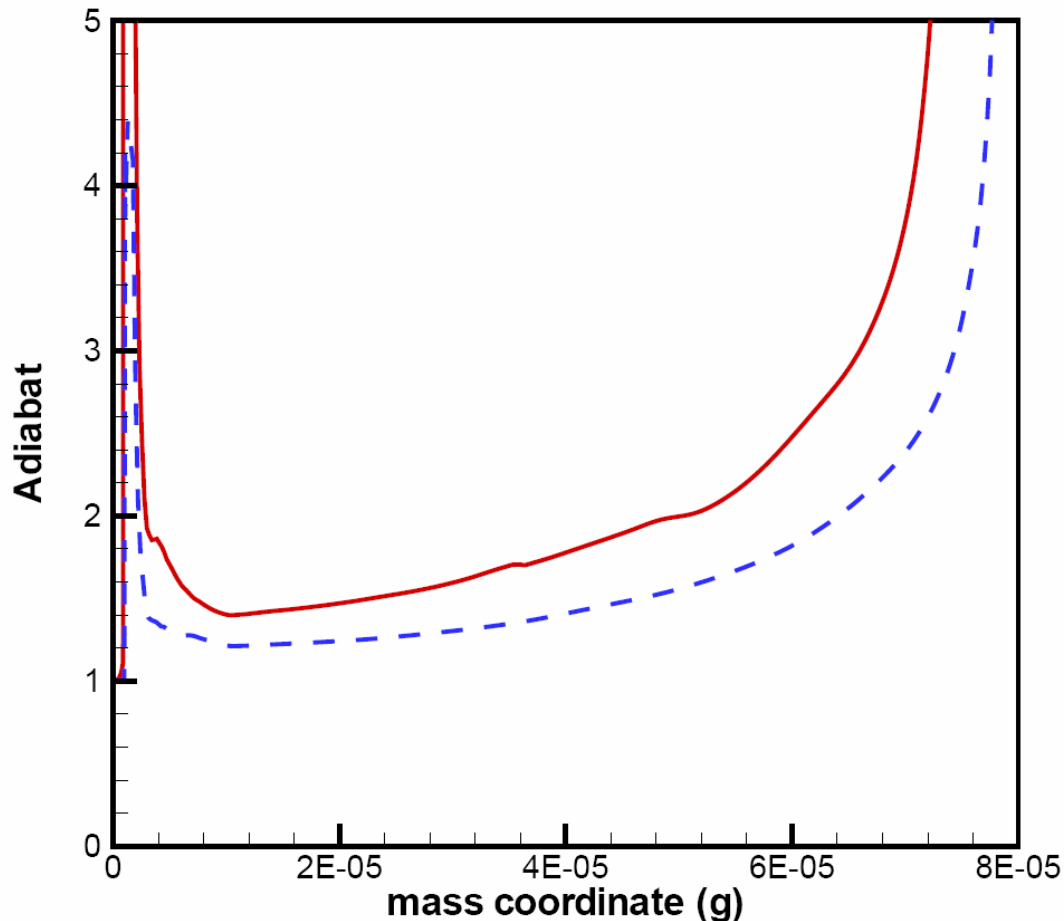
## DHe3 fill



## DD fill



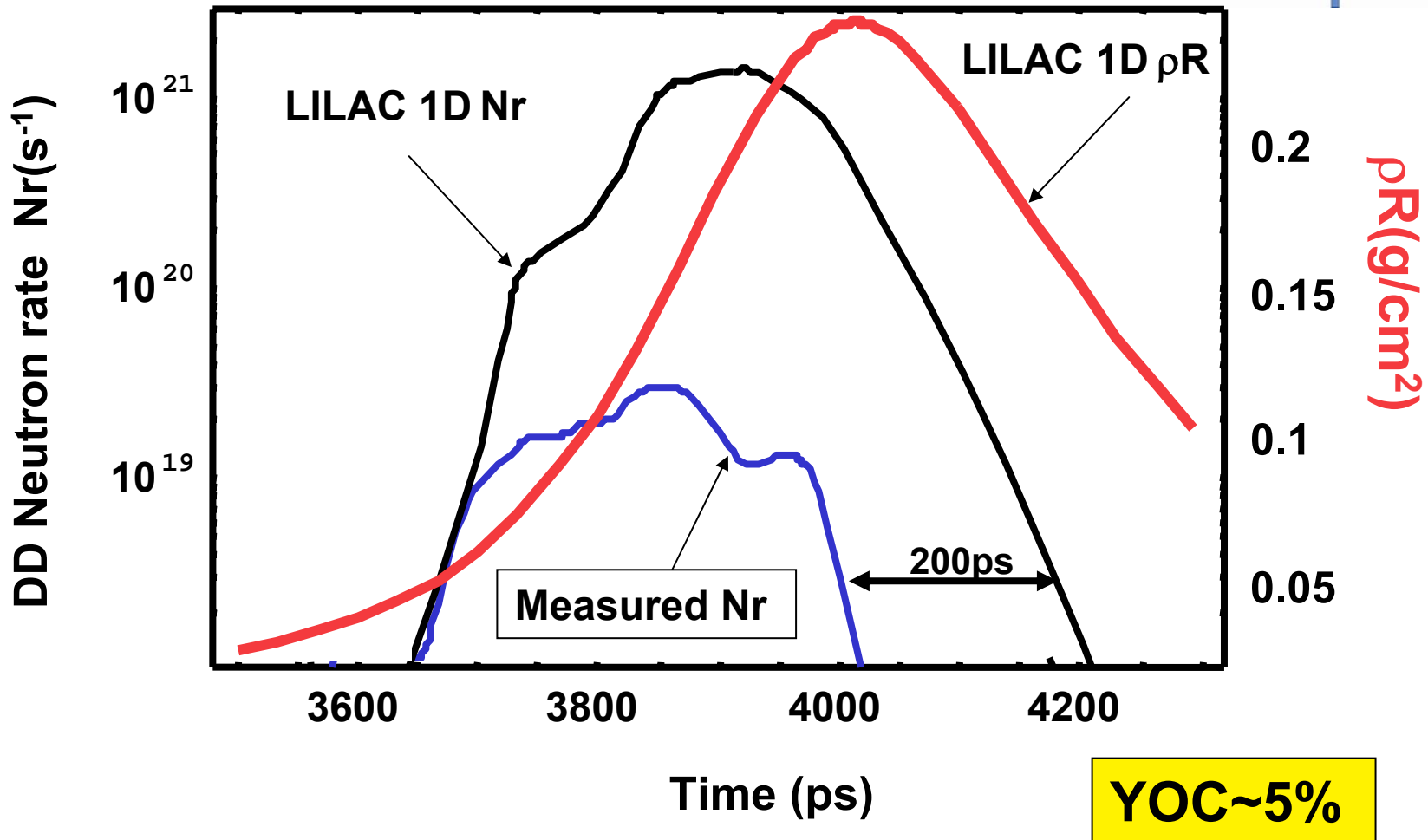
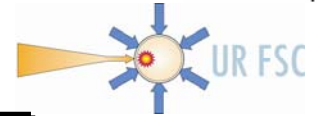
# The shell adiabat is expected to be close to the design value



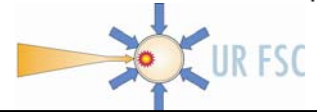
  
Adiabat simulated  
with the experimental  
laser pulse

  
Adiabat simulated  
with the design  
laser pulse

The DD burn begins as predicted and shows a 200ps truncation probably due to hot spot CH-DD mixing



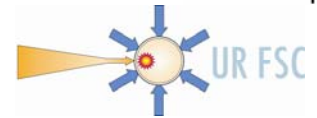
For the DD shots, the measured  $\rho R$  is compared with the simulated  $\rho R$  averaged over the measured neutron rate.



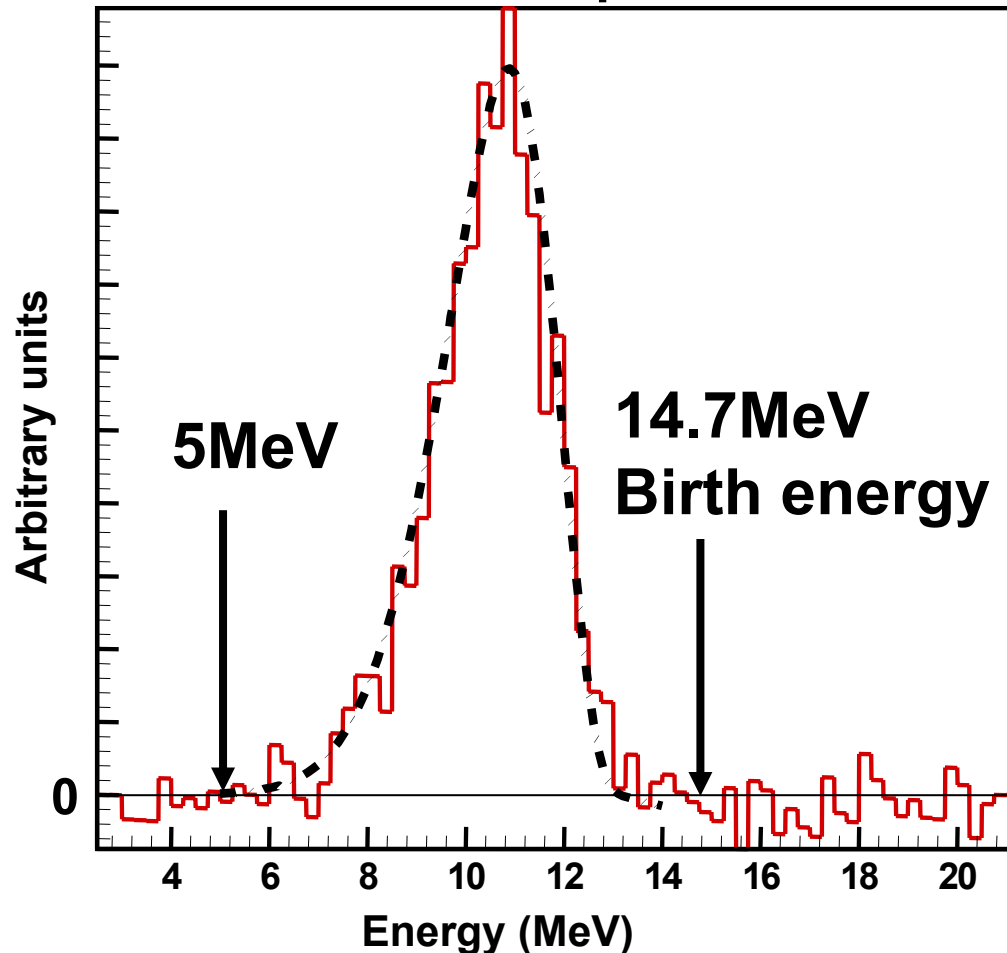
Shot Number	Gas fill	Pressure (atm)	Measured Burn-averaged $\rho R$ (g/cm <sup>2</sup> )	Simulated Burn-averaged $\rho R$ (g/cm <sup>2</sup> )
43074	D2	34	0.133	0.138
43075	D2	25	0.146	0.144
43107	D2	25	0.122	0.132
43114	D2	25	0.128	0.112
43106	D2	13	0.128	No NTD avail
43108	D2	13	0.129	No NTD avail
43109 + 43112	D3He	33	0.128	No NTD avail
43113 + 43113	D3He	25	0.130	No NTD avail
<b>AVERAGE</b>			<b>0.131</b>	<b>0.130</b>



The maximum  $\rho R$  during the burn can be inferred from the downshift of the tail of the proton spectrum



D He3 Proton spectrum



$$\begin{aligned}
 &9.7 \text{ MeV } \Delta E \\
 &- \\
 &1.5 \text{ MeV } * \\
 &\text{broadening} \\
 &= \\
 &8.2 \text{ MeV} \\
 &\text{downshift}
 \end{aligned}$$

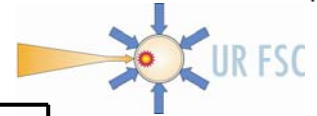


$$(\rho R)_{\text{max}}^{\text{burn}} \approx 0.25 \text{ g/cm}^2$$

\* V. Smalyuk et al, Phy. Rev. Lett. 90, 135002-1, (2003)

# The measured maximum $\rho R$ during the burn agrees with the 1D predictions

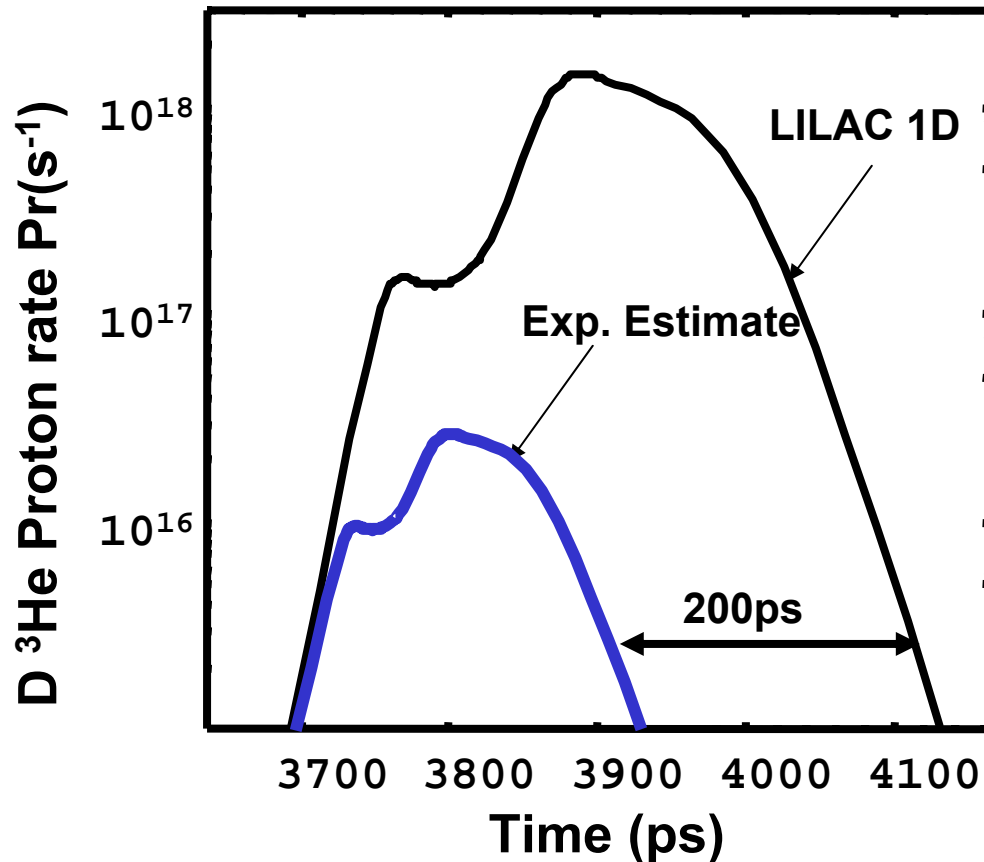
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Maximum $\rho R$ during burn	Measured (g/cm <sup>2</sup> )	Simulation (g/cm <sup>2</sup> )
$\rho R_{max}$ DHe <sub>3</sub> fill, 33 atm. Averaged over 43109-43112	0.238	0.246
$\rho R_{max}$ DHe <sub>3</sub> fill, 25 atm. Averaged over 43110-43113	0.244	0.264

- This assembly can stop 3MeV electrons for  $2\rho R$ .
- Simulations show that empty shells reach 0.5g/cm<sup>2</sup> and can stop up to 5MeV electrons for  $2\rho R$ .

The D-He3 proton rate is estimated based on the rise time and the measured 200ps truncation



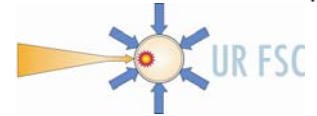
$A$  from  $Y_p = 3\%$  of 1D

$t_0$ (ns) from 1D rise time

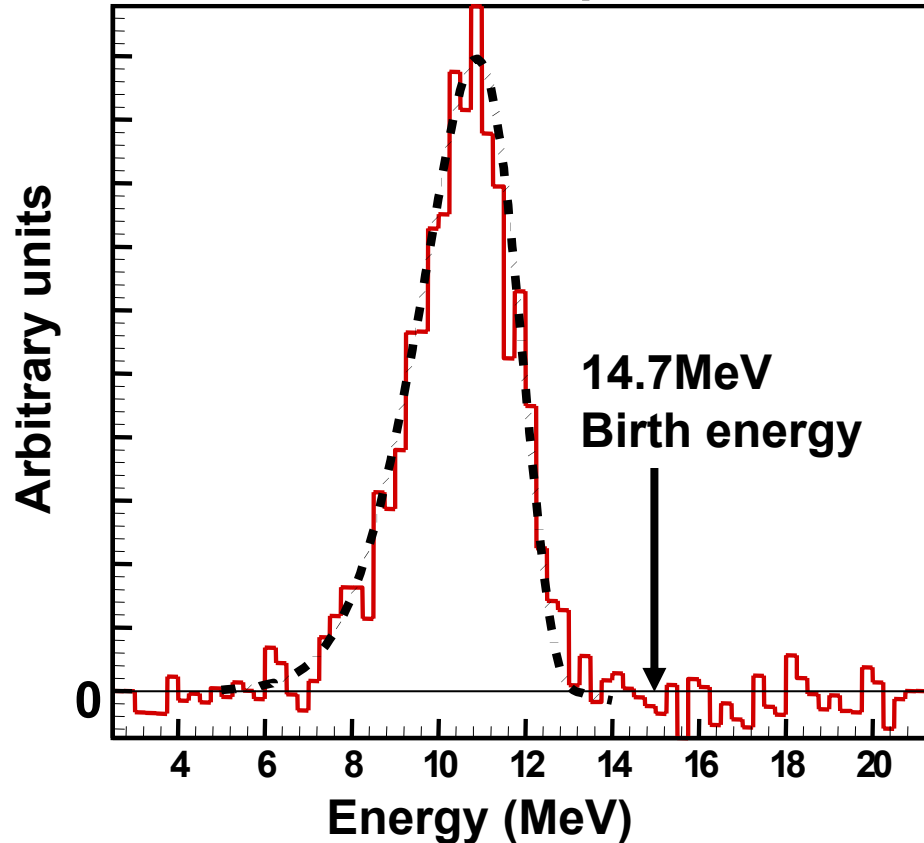
$\xi$  from 200ps burn truncation

$$\Pr_{\text{exp}}^{\text{estimate}} = A \cdot \Pr_{1D}^{\text{LILAC}} \left[ \frac{t - t_0 - 3.9}{\xi} - 3.9 \right]$$

Using the estimated proton rate and the simulated  $\rho R$ , the downshifted proton spectrum is constructed and compared with the measured spectrum



DHe3 Proton spectrum



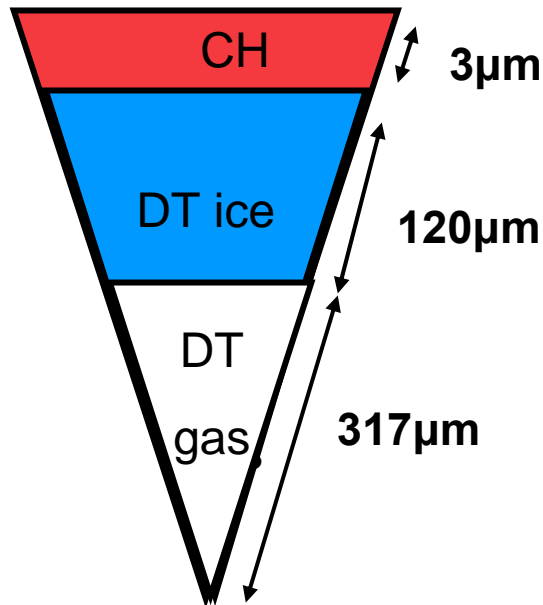
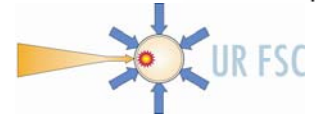
.....  
Spectrum reconstructed  
from 1D simulations  
(includes broadening of  
1.5 MeV at 5 MeV)

—————  
Measured spectrum

The good agreement suggest that the simulated time evolution of the areal density is correct

# High- $\rho_R$ Fast-Ignition 20kJ OMEGA cryo-DT implosions are designed with a low velocity and low adiabat ( $\sim 1$ )

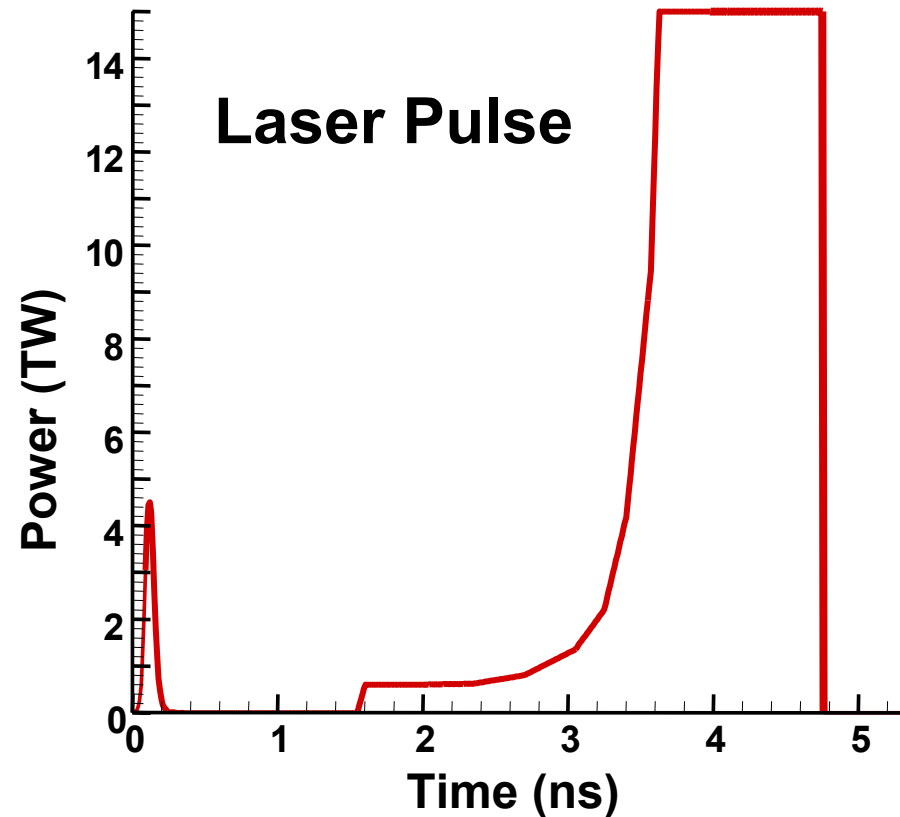
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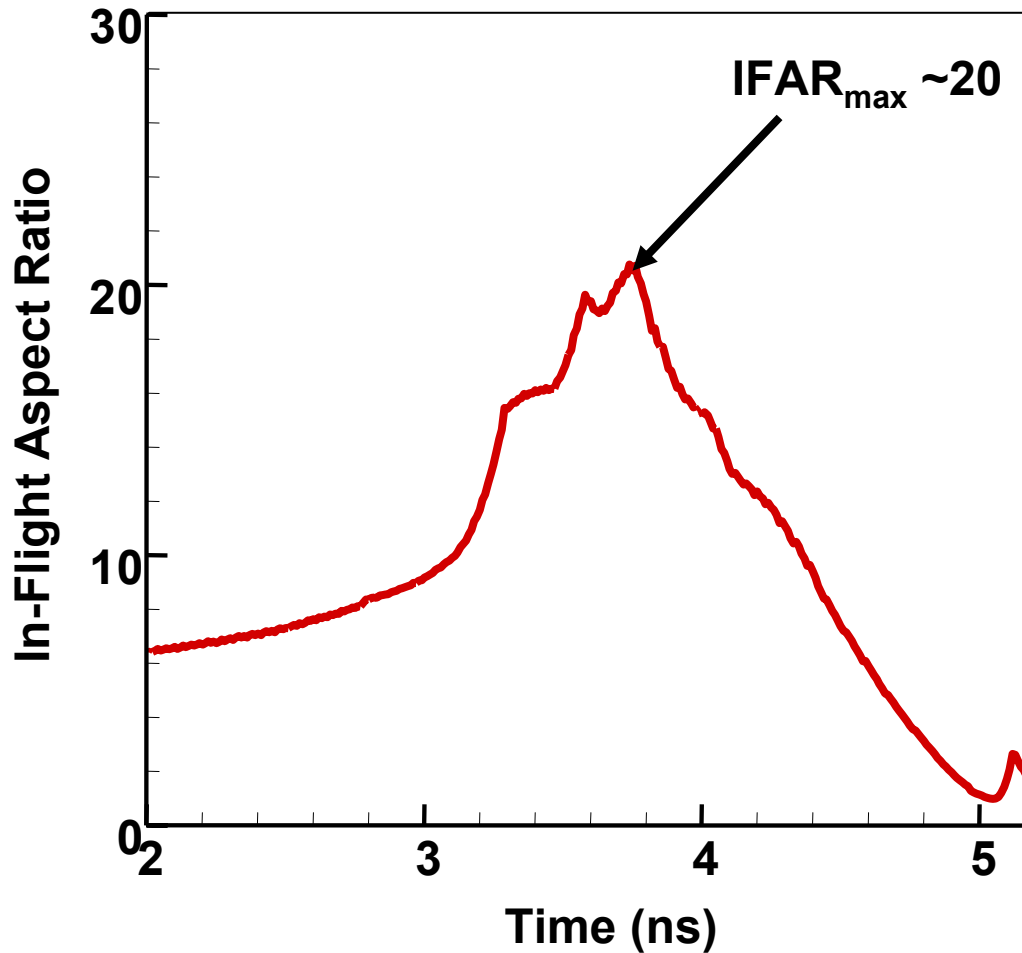
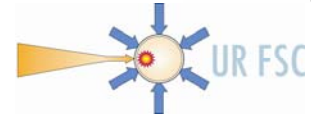
$$E_L = 20\text{kJ}$$

$$\alpha = 1.3$$

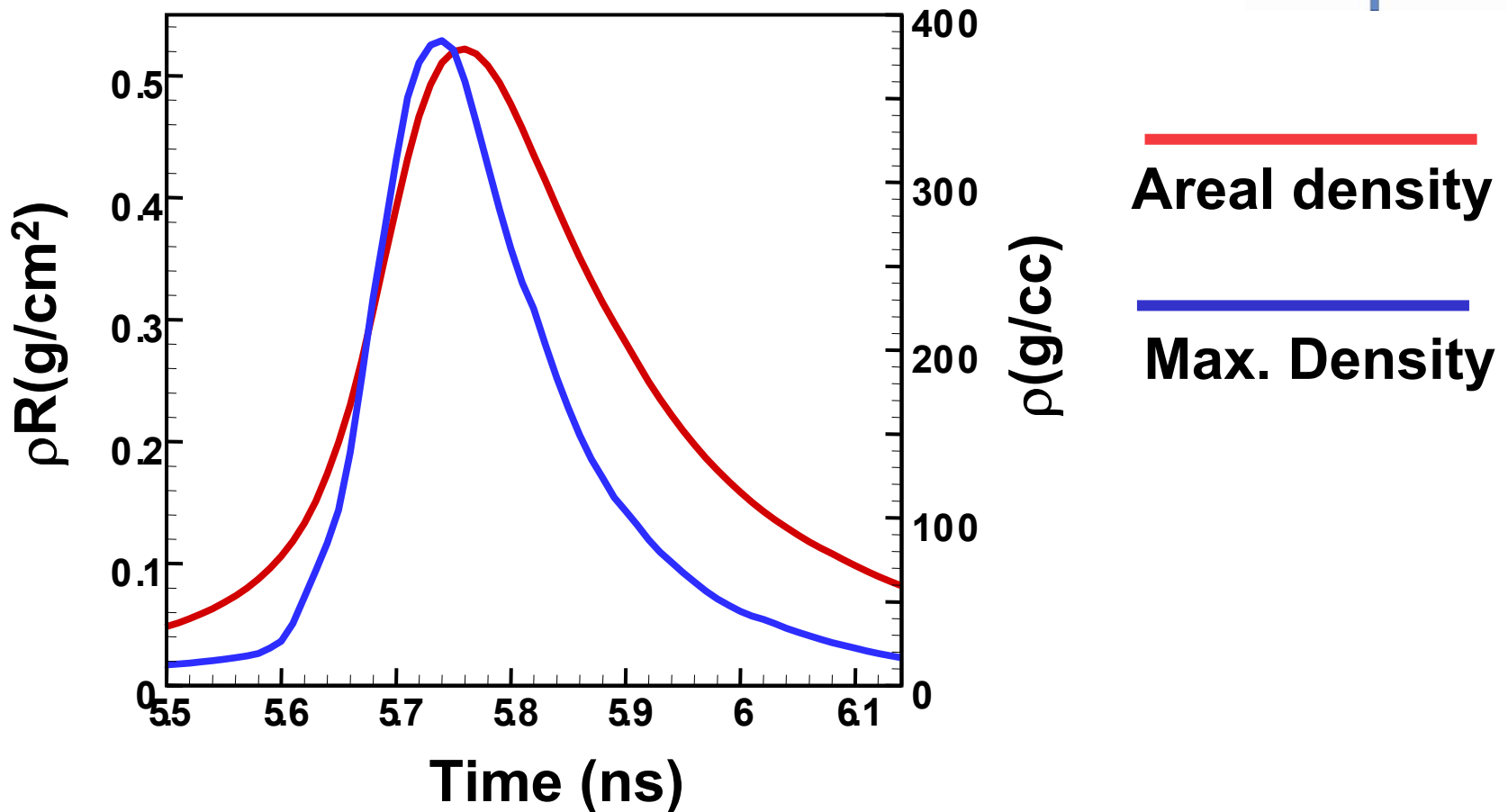
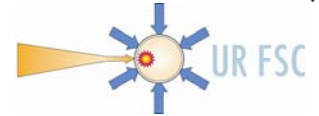
$$N\text{-yield} = 7e12$$



# The in-flight aspect ratio is low $\sim 20$ leading to 3 e-folding of RT growth

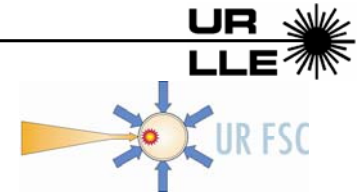


The simulated areal density exceeds  $0.5 \text{ g/cm}^2$  and the density reaches 1700 X solid density



# Fuel assembly experiments on OMEGA indicate that high areal densities can be reached with slow low-adiabat implosions

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- A gain formula for FI based on constrained optimization of the target design is derived and compared with 2D simulations of the ignition and burn
- Neutron averaged areal densities of 130-140mg/cm<sup>2</sup> have been measured from the average downshift of the primary and secondary proton spectra
- The maximum areal density during the burn inferred from the tail of the primary proton spectra reaches ~240mg/cm<sup>2</sup> in agreement with 1D simulations
- OMEGA size FI cryo targets are designed to reach areal densities above 500mg/cm<sup>2</sup> and peak densities of 380g/cc